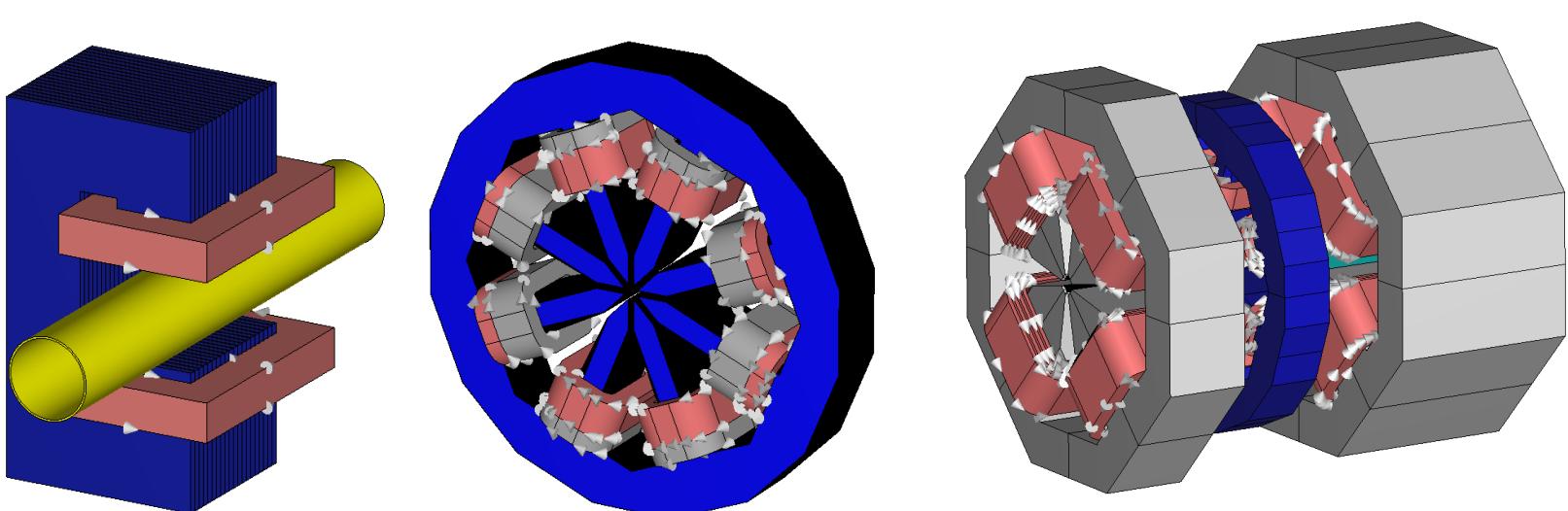
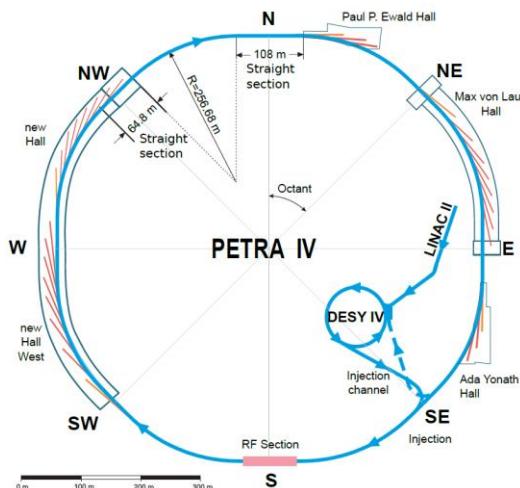


FINITE ELEMENT SIMULATION OF FAST CORRECTOR MAGNETS FOR PETRA IV

Jan-Magnus Christmann¹, Moritz von Tresckow¹, Herbert De Gersem¹,
Alexander Aloev², Sajjad H. Mirza², Sven Pfeiffer², and Holger Schlarb²



¹TEMF, TU Darmstadt, Germany

²DESY, Hamburg, Germany



CONTENTS

1 Introduction

2 Homogenization Technique

3 Toy Model

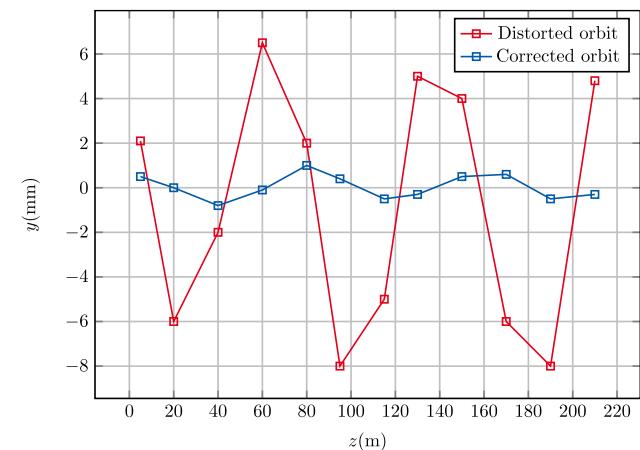
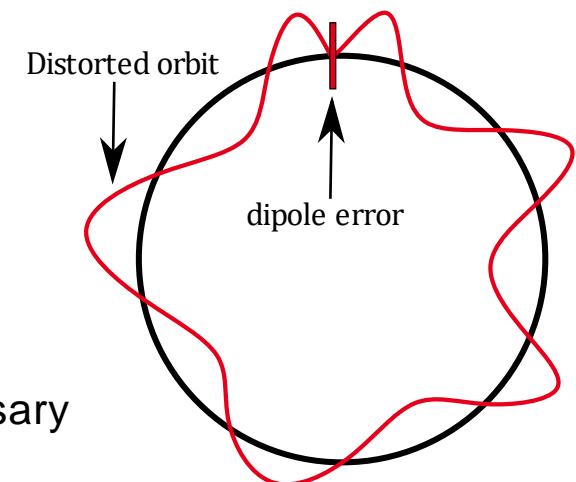
4 Stand-Alone Corrector Magnet

5 Corrector Magnet with Neighboring
Quadrupoles

6 Conclusion/Outlook

INTRODUCTION

- Circular accelerators need dipole magnets to correct orbit distortions
- **PETRA IV**: ultra-low emittance synchrotron radiation source
 - fast orbit feedback system, **corrector magnets with frequencies in kHz range** necessary
- **Strong eddy currents** → power losses, time delay, and field distortion
- **Simulation challenging** due to small skin depths and laminated yoke
 - Need for technique to simplify simulations





CONTENTS

1 Introduction

2 Homogenization Technique

3 Toy Model

4 Stand-Alone Corrector Magnet

5 Corrector Magnet with Neighboring
Quadrupoles

6 Conclusion/Outlook

THEORY

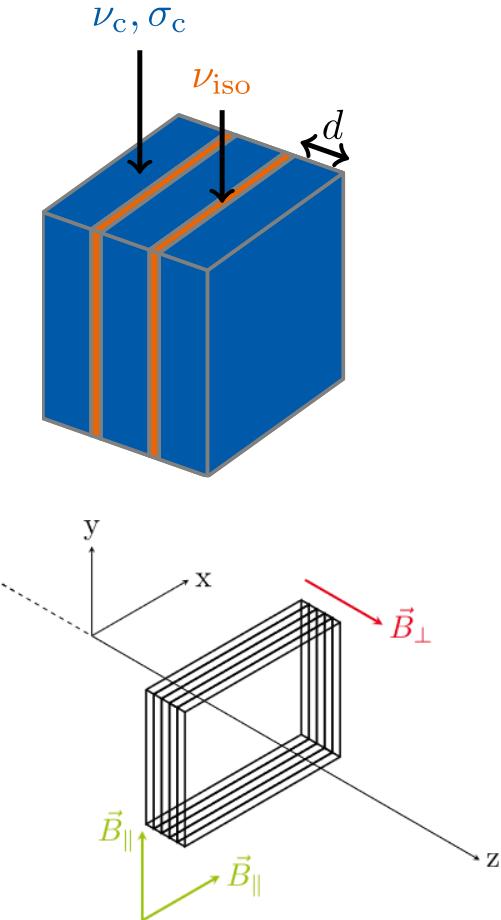
- Magnetoquasistatic PDE: $\nabla \times (\nu(\vec{r}) \nabla \times \underline{\vec{A}}(\vec{r})) + j\omega\sigma(\vec{r})\underline{\vec{A}}(\vec{r}) = \vec{J}_s(\vec{r})$
- Replace reluctivity $\nu(\vec{r})$ and conductivity $\sigma(\vec{r})$ in the laminated yoke with spatially constant tensors

$$\nu(\vec{r}) \rightarrow \bar{\nu} = \frac{1}{8}\sigma_c d \delta \omega (1+j) \frac{\sinh((1+j)\delta^{-1}d)}{\sinh^2((1+j)\delta^{-1}d/2)} \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \end{bmatrix} + \nu_c \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

$$\sigma(\vec{r}) \rightarrow \bar{\sigma} = \gamma \sigma_c \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

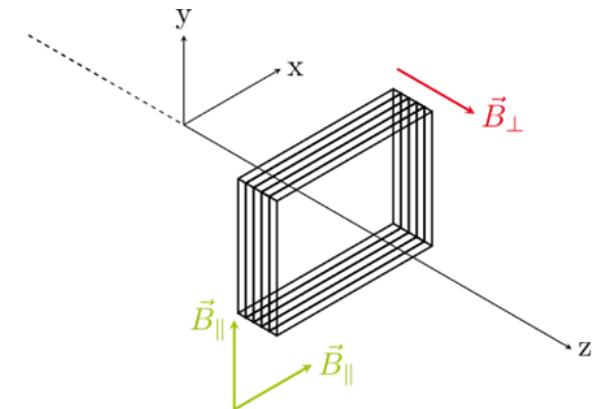
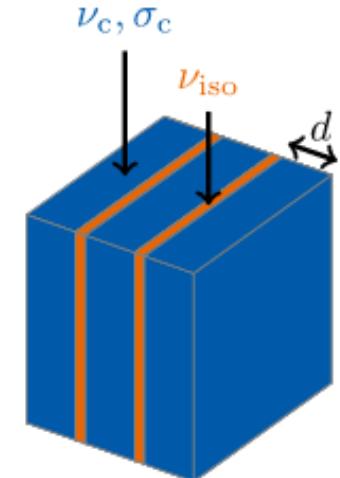
Skin depth $\delta = \sqrt{2/\omega\sigma_c\mu_c}$
 Stacking factor $\gamma = \frac{V_c}{V_{Yoke}}$

P. Dular et al., 2003
 L. Krählenbühl et al., 2004
 H. De Gersem et al., 2012



APPLICATION

- Frequency-dependent, complex-valued and anisotropic materials **can be implemented in LF frequency domain solver of CST Studio Suite®**
- Homogenization captures losses due to eddy currents induced by in-plane and perpendicular flux components
- Homogenization is valid also for high frequencies, i.e., $\delta \ll d$
- Restriction to frequency domain simulations
- Non-linear material properties and hysteresis are neglected





CONTENTS

1 Introduction

2 Homogenization Technique

3 Toy Model

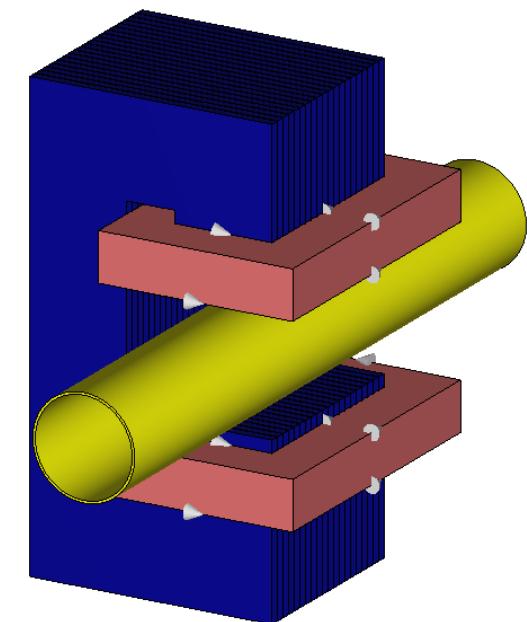
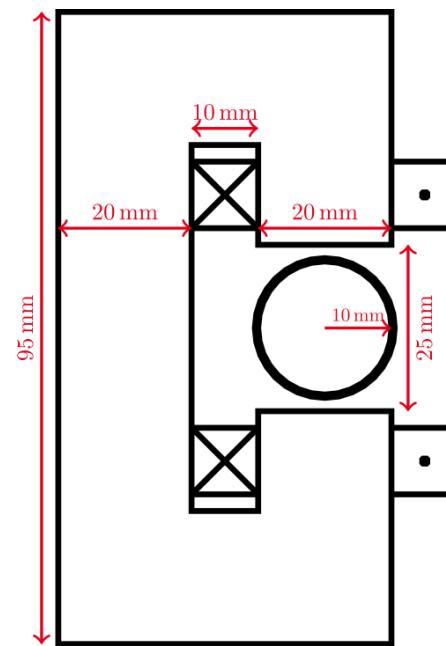
4 Stand-Alone Corrector Magnet

5 Corrector Magnet with Neighboring
Quadrupoles

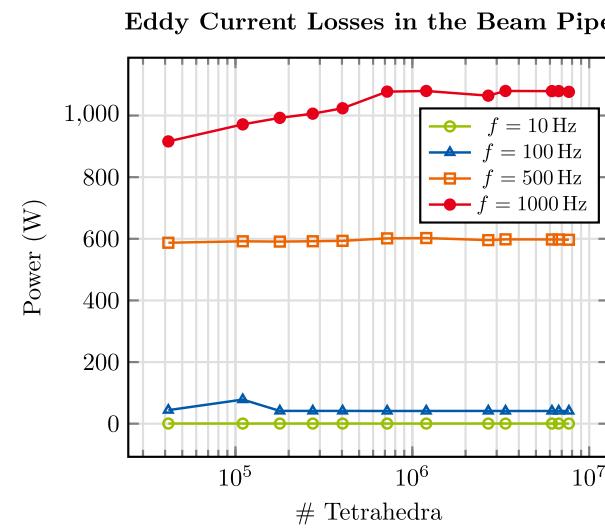
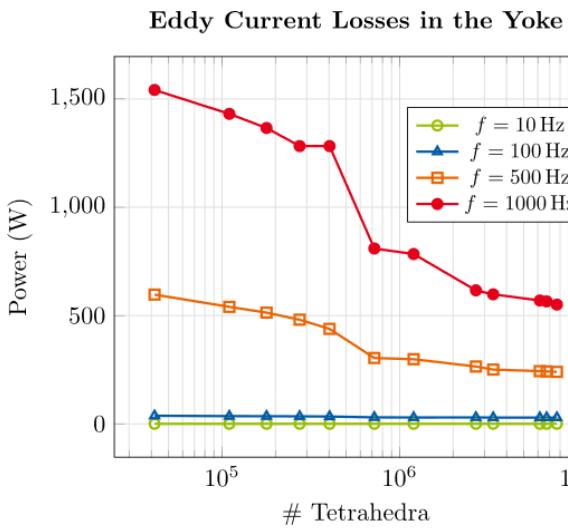
6 Conclusion/Outlook

MODEL DESCRIPTION

- **Iron yoke:** length = 40 mm, lamination thickness = 1.83 mm
- **Copper beam pipe:** thickness = 0.5 mm, length = 140 mm
- **Coils:** current = 10 A (peak), # turns = 250
- **Frequency domain simulation** via **CST Studio Suite®**



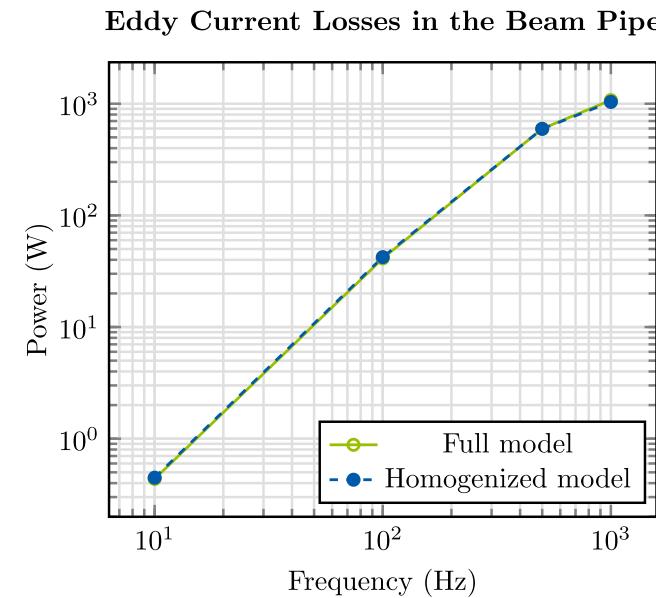
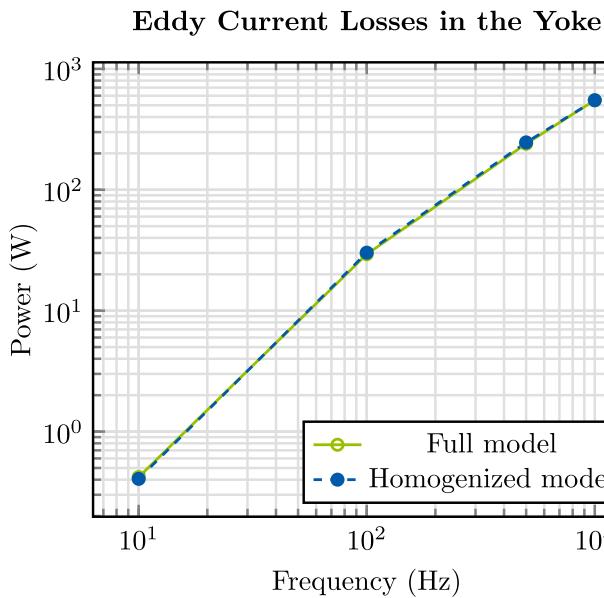
SIMULATION OF THE FULL MODEL



- Strong mesh dependence of power losses at higher frequencies
 - Obtaining reliable results is difficult
 - Need for simplified model

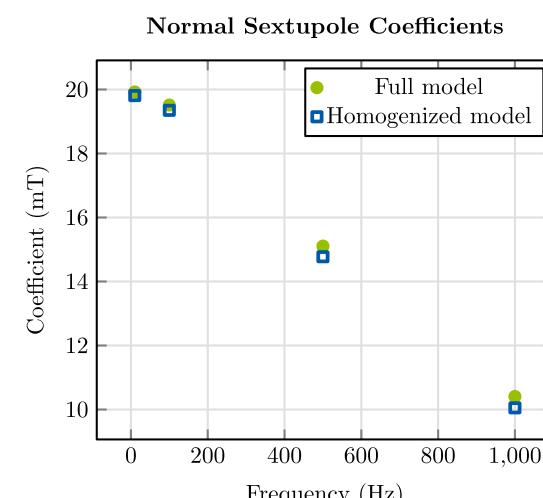
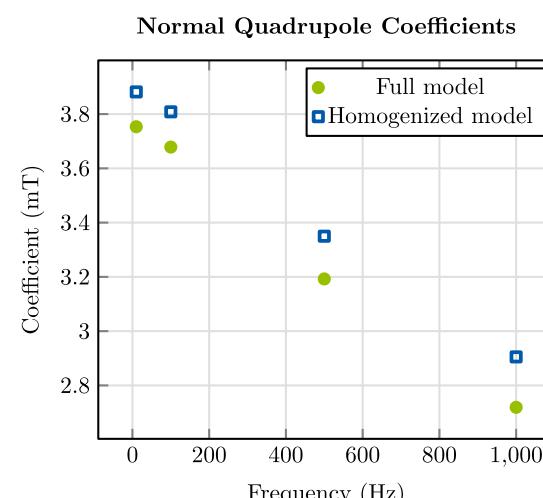
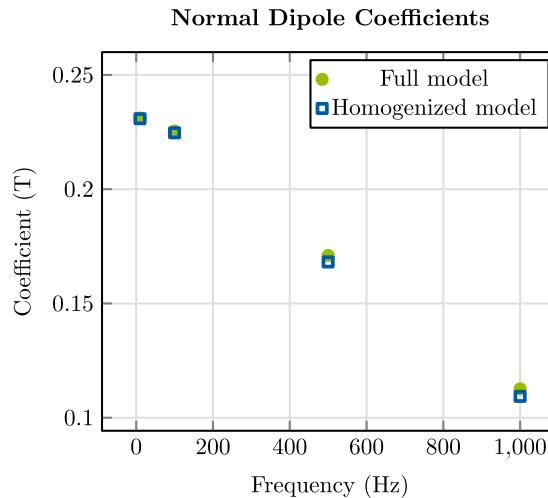
# Tetrahedra	$4.2 \cdot 10^4$	$4.0 \cdot 10^5$	$1.2 \cdot 10^6$	$3.4 \cdot 10^6$	$7.7 \cdot 10^6$
Simulation time	2 min	20 min	1 h	7.5 h	21.5 h

HOMOGENIZED VS. FULL MODEL



- Good approximation of losses in yoke & beam pipe (max. relative error 4 %)
- Simulation time reduced from several hours to 4 min

HOMOGENIZED VS. FULL MODEL



- Homogenization technique yields accurate approximation of multipole coefficients
→ Aperture field accurately represented

Multipole coefficient	Average rel. error
Dipole	1 %
Quadrupole	5 %
Sextupole	2 %



CONTENTS

1 Introduction

2 Homogenization Technique

3 Toy Model

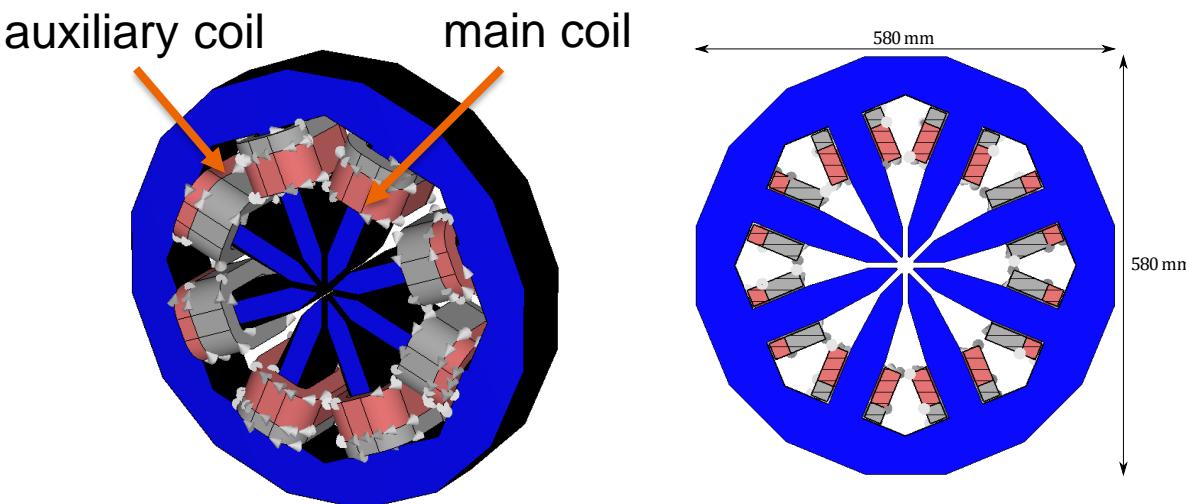
4 Stand-Alone Corrector Magnet

5 Corrector Magnet with Neighboring
Quadrupoles

6 Conclusion/Outlook

MODEL DESCRIPTION

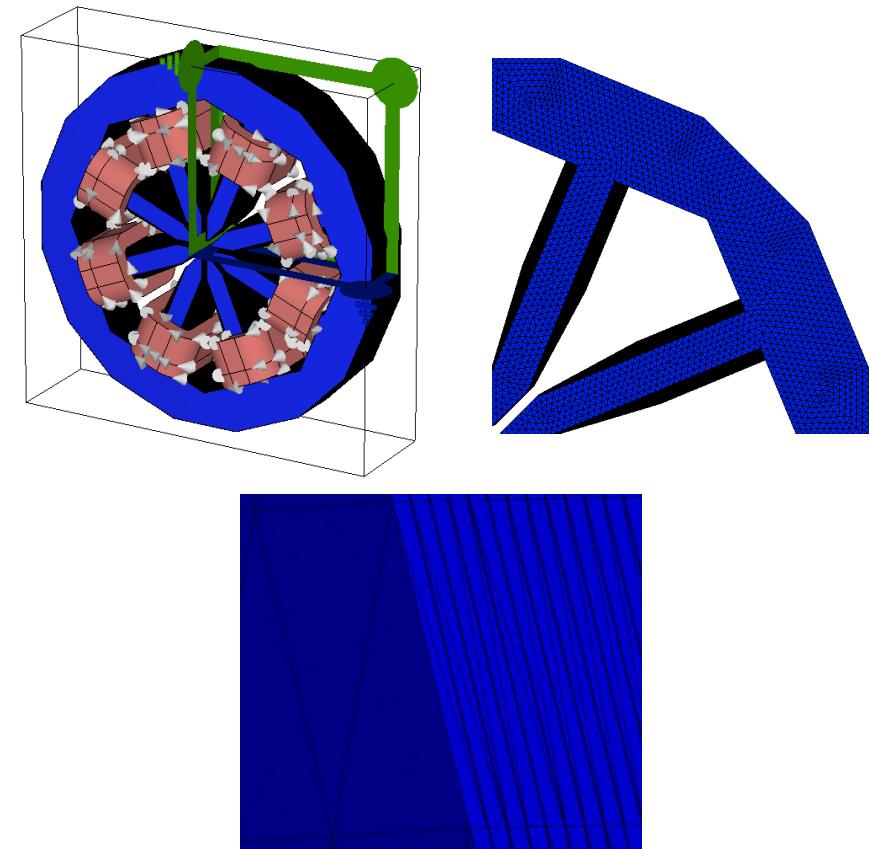
- **Dipole** corrector with **octupole-like design**
- **Coils:**
 - 4 main coils: current = 27.4 A (peak), # turns = 53
 - 4 auxiliary coils: current = 27.4 A (peak), # turns = 22
- **Iron yoke:**
 - Diameter = 580 mm, length = 90 mm
 - Lamination thickness = 0.5 mm
- At first **no beam pipe**



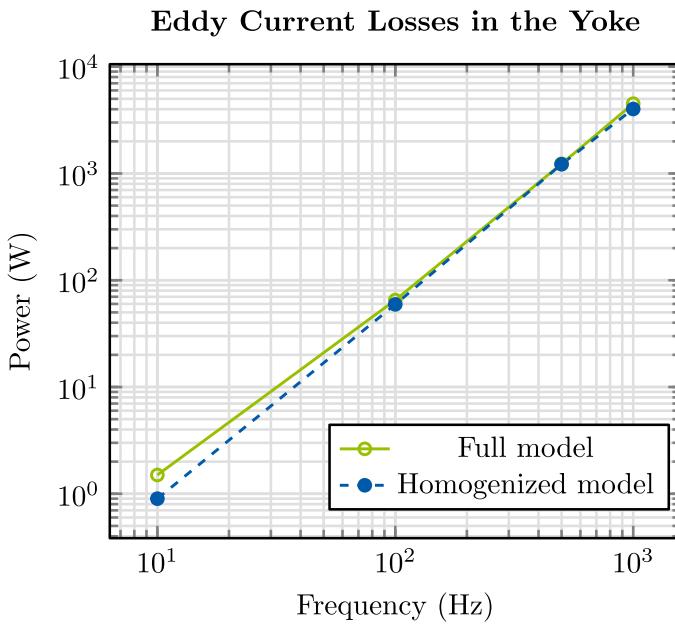
Design by A. Aloev (DESY),
inspired by APS

SIMULATION OF THE FULL MODEL

- Frequency domain simulation via **CST Studio Suite®**
- Three symmetry planes, test frequencies $f = 10 \text{ Hz}, 100 \text{ Hz}, 500 \text{ Hz}, 1 \text{ kHz}$
- Long simulation times even for relatively coarse meshes
- Finest mesh: # tetrahedra = $2.3 \cdot 10^6 \rightarrow$ simulation time = 26 h
- Skin depth cannot be resolved \rightarrow power loss still mesh-dependent



HOMOGENIZED VS. FULL MODEL

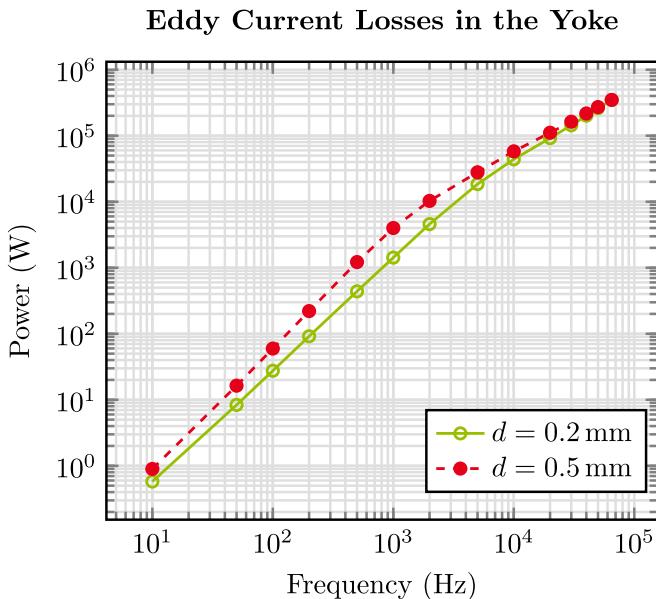


Multipole coefficient	Average rel. deviation
Dipole	1 %
14-pole	1 %
18-pole	3 %

- Similar power losses
 - Good agreement in multipole coefficients
 - Simulation time reduces from 26 h to 5 min
- Homogenized model can be used for further studies

Keep in mind:
Power losses in full model
are still mesh-dependent !

LOSSES FOR DIFFERENT LAMINATION THICKNESSES

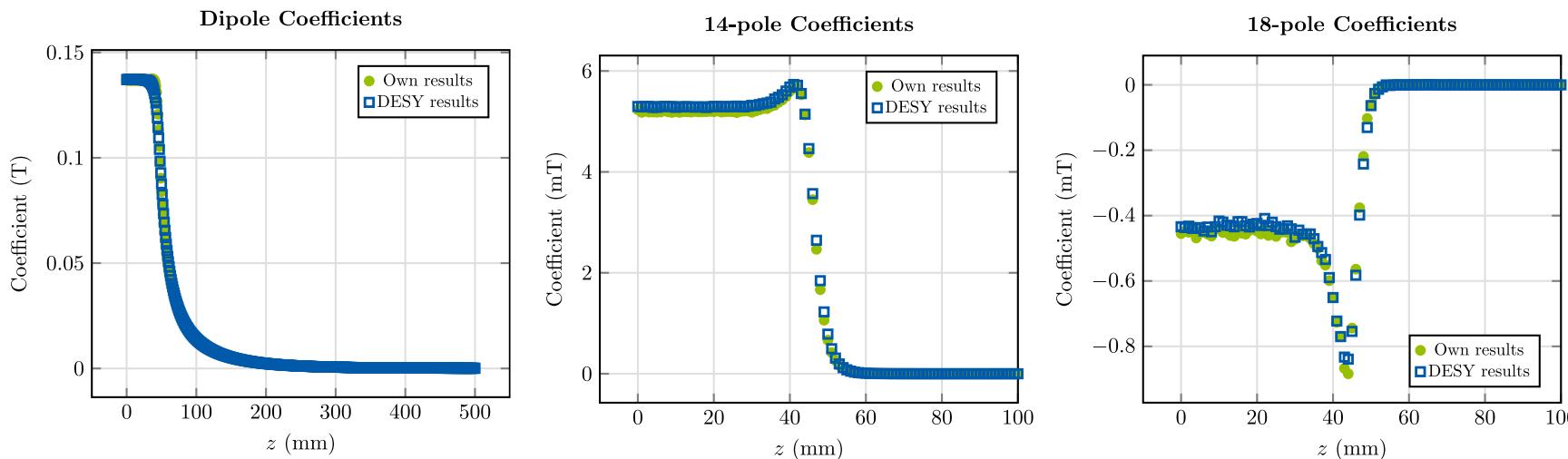


f (Hz)	Eddy current losses (W)			
	$d = 0.2$ mm	$d = 0.3$ mm	$d = 0.4$ mm	$d = 0.5$ mm
10	$5.8 \cdot 10^{-1}$	$6.5 \cdot 10^{-1}$	$7.6 \cdot 10^{-1}$	$9.0 \cdot 10^{-1}$
100	$2.8 \cdot 10^1$	$3.4 \cdot 10^1$	$4.6 \cdot 10^1$	$6.0 \cdot 10^1$
500	$4.4 \cdot 10^2$	$6.2 \cdot 10^2$	$9.0 \cdot 10^2$	$1.2 \cdot 10^3$
1000	$1.4 \cdot 10^3$	$2.1 \cdot 10^3$	$3.1 \cdot 10^3$	$4.0 \cdot 10^3$
10000	$4.4 \cdot 10^4$	$4.9 \cdot 10^4$	$5.5 \cdot 10^4$	$5.8 \cdot 10^4$
30000	$1.4 \cdot 10^5$	$1.6 \cdot 10^5$	$1.6 \cdot 10^5$	$1.6 \cdot 10^5$
65000	$3.5 \cdot 10^5$	$3.6 \cdot 10^5$	$3.6 \cdot 10^5$	$3.5 \cdot 10^5$

Simulation uses the same current for all frequencies !

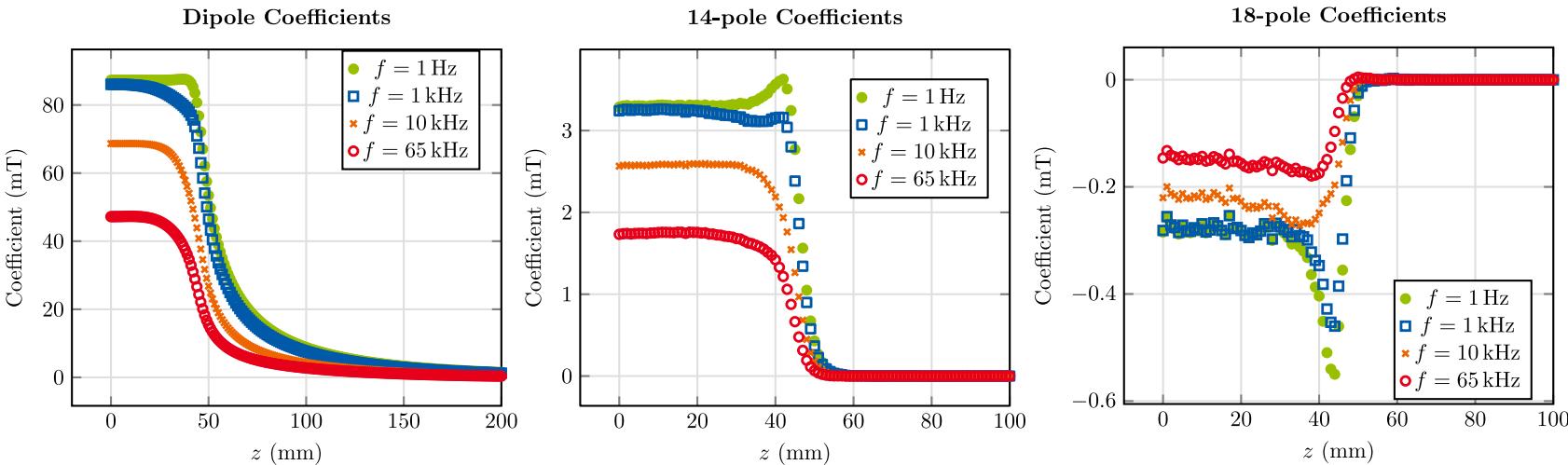
- Use homogenization to investigate losses up to 65 kHz
- Vary $d = 0.2 - 0.5$ mm, keep $\gamma \approx 0.91$ constant
- At **low frequencies**, the lamination thickness has **strong influence** on the losses
- At **very high frequencies**, the lamination thickness has **no influence** on the losses

LONGITUDINAL MULTIPOLE DISTRIBUTION (STATIC)



- Compute multipole coefficients along longitudinal axis of the magnet
- Comparison with DESY for static case → **good agreement**

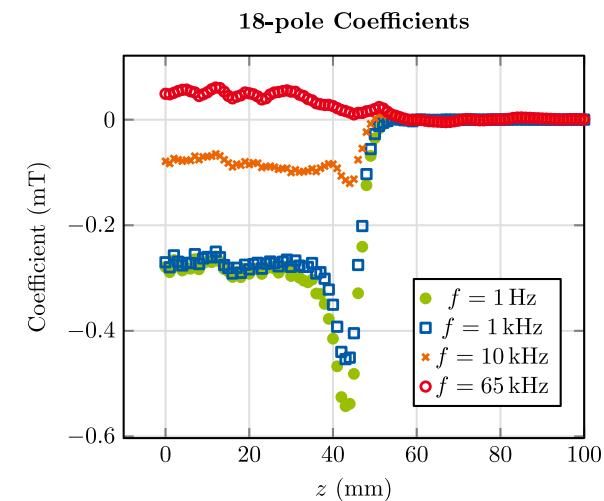
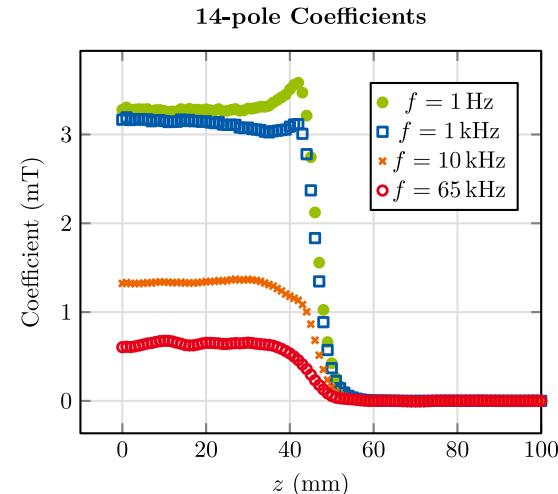
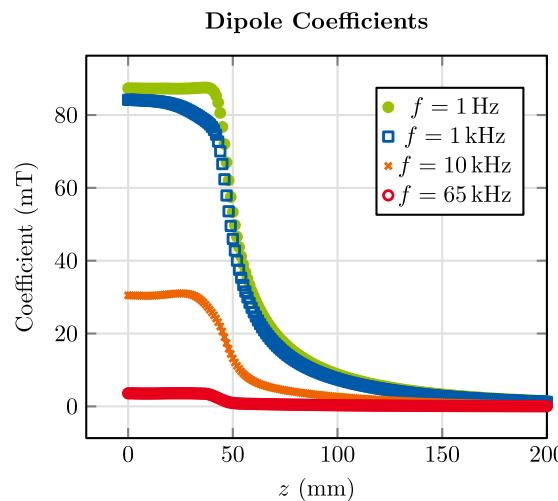
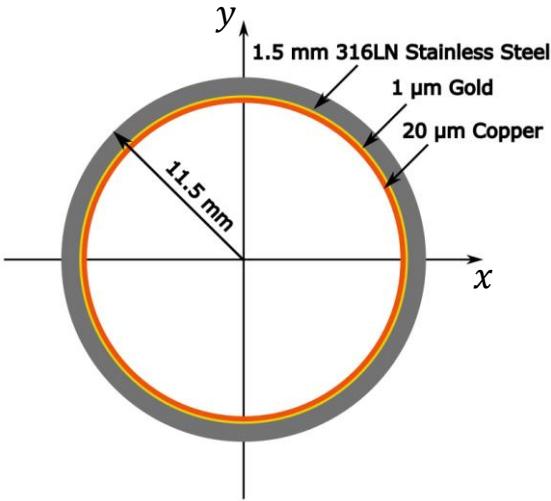
LONGITUDINAL MULTIPOLE DISTRIBUTION (TIME-HARMONIC)



- **Updated # turns & current:**
 - Main coils: 65 turns, 15 A
 - Aux. coils: 27 turns, 15 A
- 65 kHz vs. 1 Hz:
 - Int. dipoles: -57 %
 - Int. 14-poles: -52 %
 - Int. 18-poles: -54 %

f (Hz)	Int. dipole (mT m)	Int. 14-pole (μ T m)	Int. 18-pole (μ T m)
1	11.6	316.4	-30.3
1000	10.7	300.4	-28.6
10000	7.6	229.0	-21.5
65000	5.0	150.3	-13.9

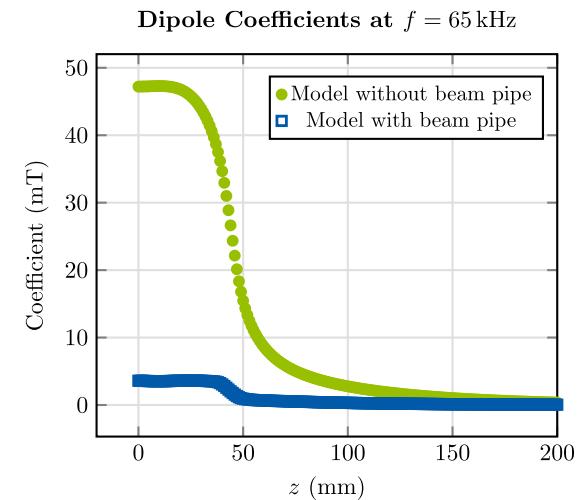
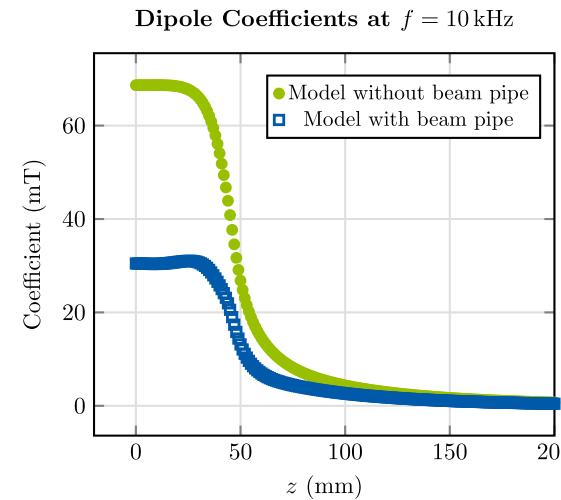
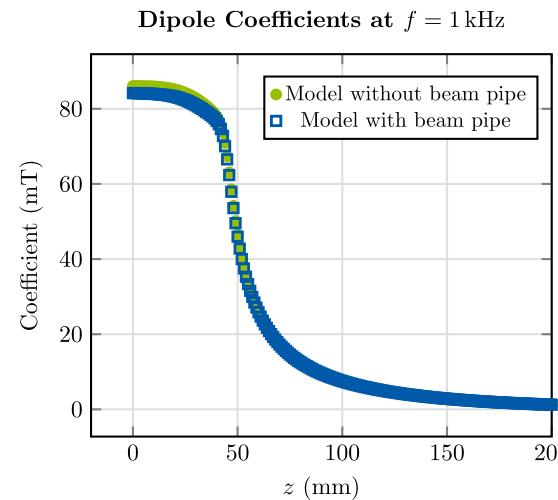
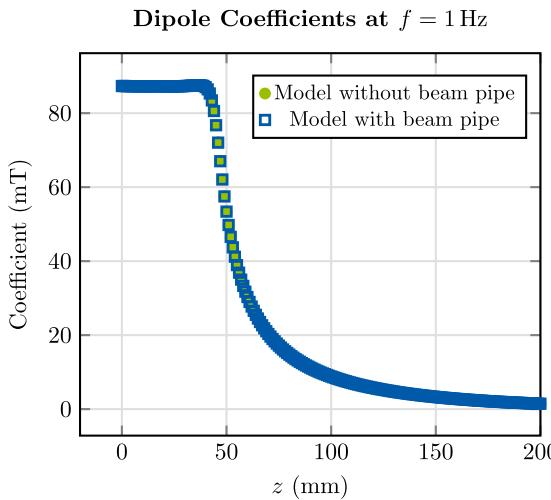
INCLUSION OF BEAM PIPE



f (Hz)	Int. dipole (mT m)	Int. 14-pole ($\mu\text{T m}$)	Int. 18-pole ($\mu\text{T m}$)
1	11.5	313.3	-30.6
1000	10.5	292.7	-28.1
10000	3.6	122.6	-8.3
65000	0.4	57.4	4.3

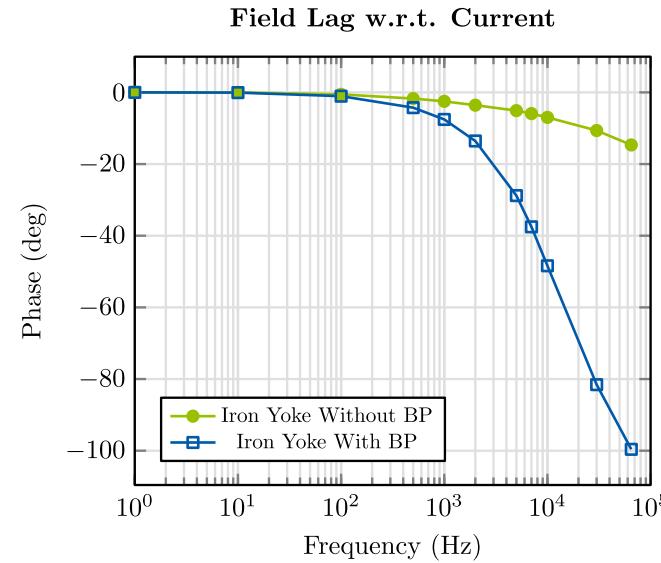
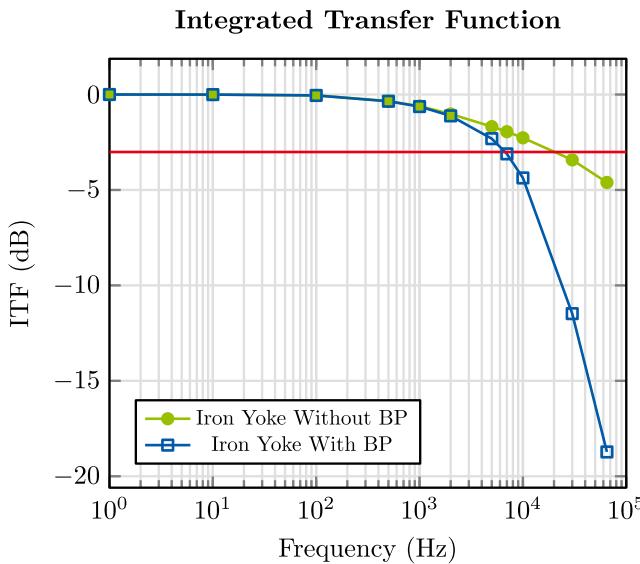
- General shape similar to model without beam pipe
- 65 kHz vs. 1 Hz:
 - Int. dipoles: -97 % (-57 %)
 - Int. 14-poles: -82 % (-52 %)
 - Int. 18-poles change sign (-54 %)

INCLUSION OF BEAM PIPE



- **Up to $f \approx 1 \text{ kHz}$:** Only minor differences between the two models
- **For $f >> 1 \text{ kHz}$:** Strong attenuation of dipole field due to eddy currents in beam pipe
- At higher frequencies, beam pipe leads to greater effective length of the magnet

INTEGRATED TRANSFER FUNCTION AND FIELD LAG



$$\text{ITF}(f) = \frac{\int_l B_1(z, f) dz}{\int_l B_1(z, f = 1\text{Hz}) dz}$$

Yoke material	3 dB bandwidth	Phase shift at bandwidth
Iron	7 kHz	38°
M-19 Steel	10 kHz	46°
1010 Steel	7 kHz	38°

Yoke material	Average relative permeability*	Conductivity (MS/m)
Iron	5690	10.4
M-19 Steel	4166	1.9
1010 Steel	2780	6.993

* Values are computed from results of static simulations with non-linear BH-curve

CONTENTS

- 1** Introduction
- 2** Homogenization Technique
- 3** Toy Model
- 4** Stand-Alone Corrector Magnet
- 5** Corrector Magnet with Neighboring Quadrupoles

- 6** Conclusion/Outlook

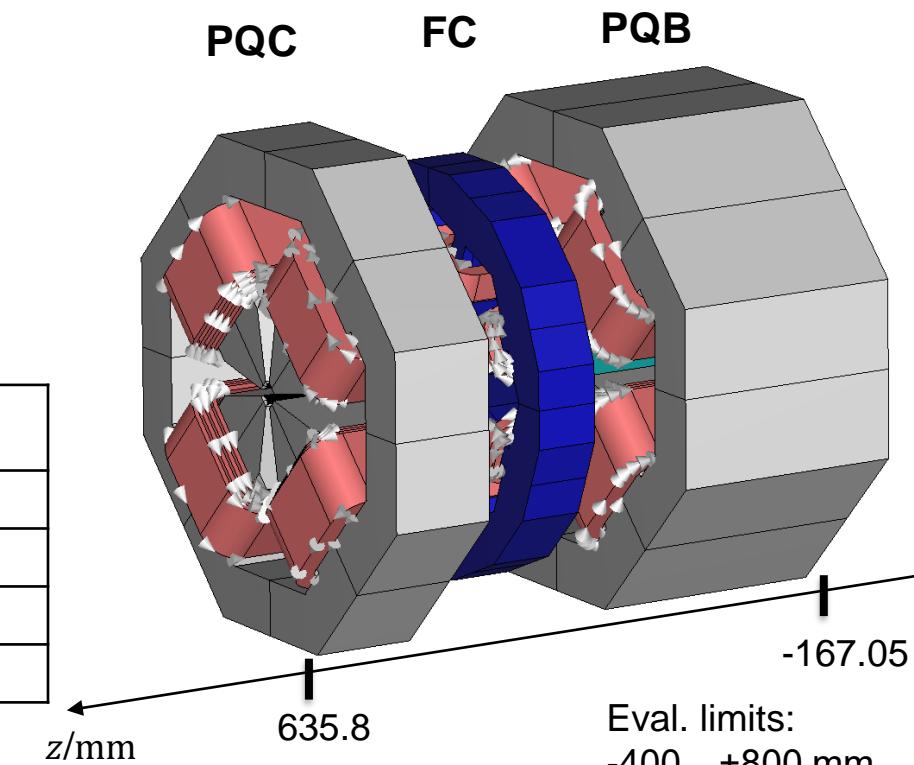
MODEL DESCRIPTION

- Corrector magnet (FC) with two neighboring quadrupole magnets (PQB & PQC)
- AC currents in corrector coils, DC currents in quadrupole coils
- All yokes are 1010 steel, PQB quadrupoles have Vacoflux-50 poles
- Quadrupole yokes are solid, corrector yoke is laminated
- Beam pipe made out of 316LN SS with outer radius of 11 mm and thickness of 1 mm
- Distance between corrector yoke and quadrupole yokes ~ 11.5 cm

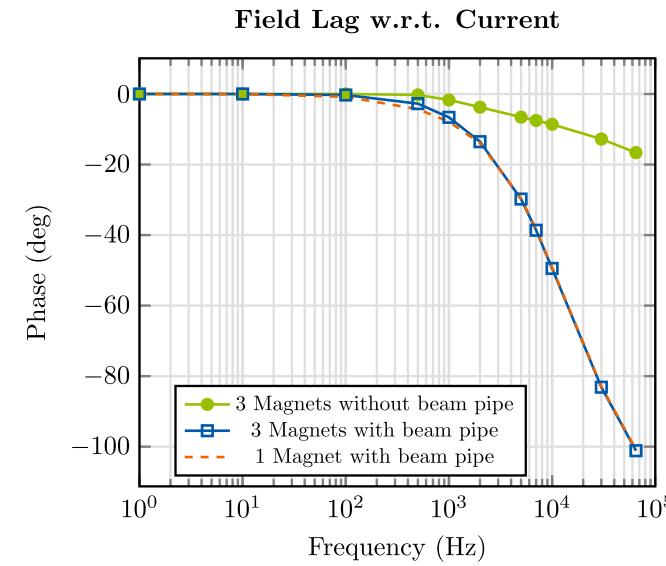
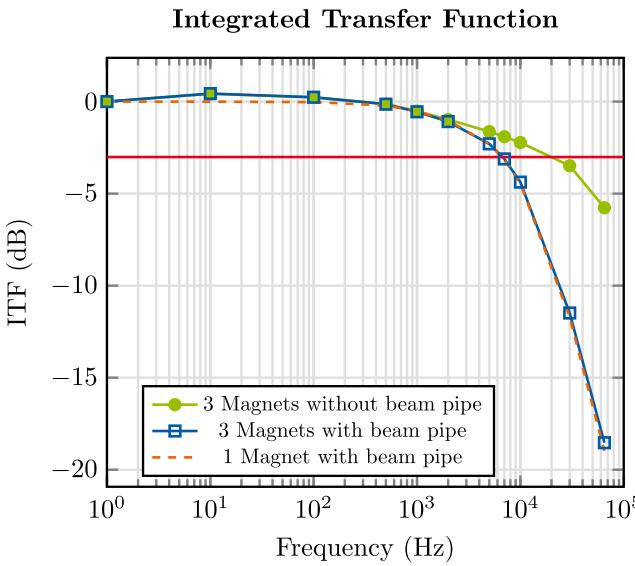
Material	Average Relative Permeability*	Conductivity (MS/m)
1010 Steel (PQC)	1450	6.993
1010 Steel (PQB)	1810	6.993
Vacfux-50 (PQB)	5000	2.38
1010 Steel (FC)	2780	6.993

Coils	Ampere turns
PQB	5728.1 At
FC (main)	975 At
FC (aux.)	405 At
PQC	5659.5 At

* Values are computed from results of static simulations with non-linear BH-curve



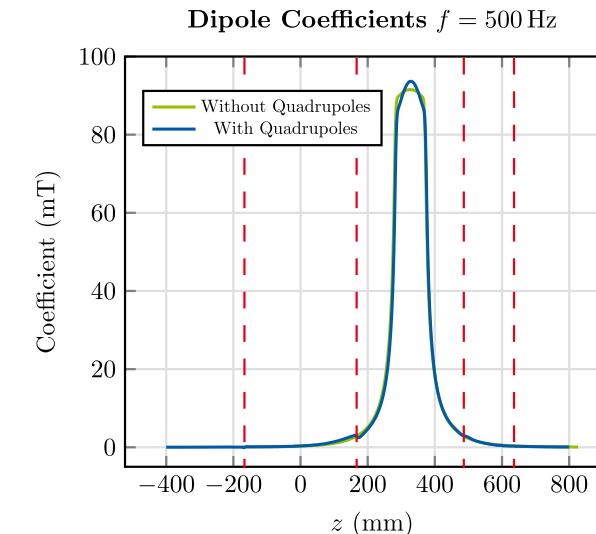
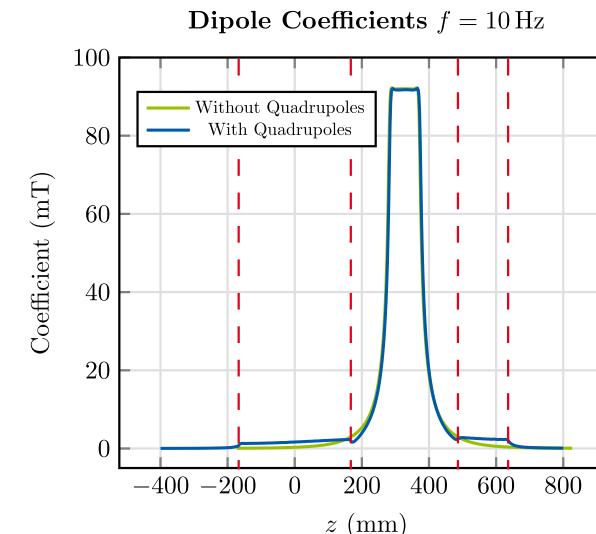
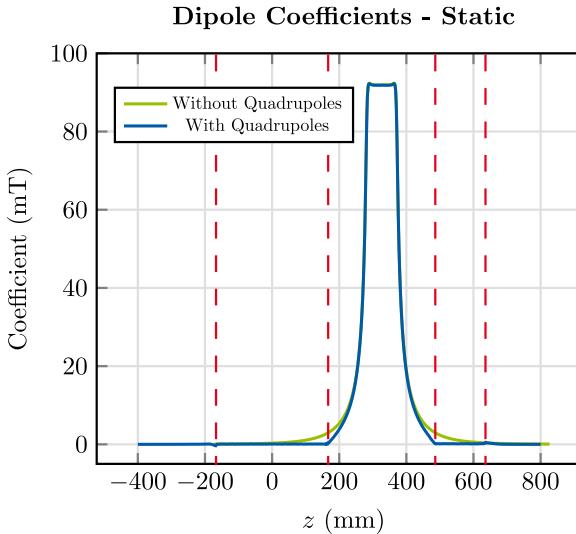
INTEGRATED TRANSFER FUNCTION AND FIELD LAG



	Model without beam pipe	Model with beam pipe
3 dB bandwidth	20 kHz	7 kHz
Phase shift at bandwidth	11°	39°

- Very similar results as for the model without neighboring quadrupoles
- **Main difference:** at low frequencies, a ~0.5 dB peak is occurring in the ITF of the model with the neighboring quadrupoles

DIPOLE COEFFICIENTS ALONG THE AXIS



Simulation with beam pipe

- At low frequencies ($f \leq 100$ Hz), we observe a **parasitic dipole** component **inside the quadrupole magnets**
 - This dipole component is due to eddy currents induced in the quadrupole yokes by the AC corrector field
- ➔ Peak in ITF at low frequencies
- ➔ Shift of the center of mass (~ 0.5 cm at most)



CONTENTS

- 1** Introduction
- 2** Homogenization Technique
- 3** Toy Model
- 4** Stand-alone Corrector Magnet
- 5** Corrector Magnet with Neighboring Quadrupoles

- 6** Conclusion/Outlook

CONCLUSION/OUTLOOK

- Validation of homogenization technique using **toy model**
 - ➔ Good approximation of multipoles and power losses
 - ➔ Simulation time reduced from several hours to a few minutes
 - Application to **corrector magnet model**
 - Power losses for different lamination thicknesses
 - Longitudinal multipole distributions
 - Integrated transfer function and field lag
 - Cross-talk with neighboring magnets
 - **Ongoing investigations:**
 - Simulations with different variations of the beam pipe and cooling channels
 - Approximate treatment of non-linear material properties
- }
- Homogenization enables us to study this over the frequency range of interest from DC up to 65 kHz

REFERENCES

- [1] PETRA IV Conceptual Design Report.
- [2] K. Wille, *Physik der Teilchenbeschleuniger und Synchrotronstrahlungsquellen*. Stuttgart, Germany: Teubner, 1992.
- [3] P. Dular et al., “A 3-D Magnetic Vector Potential Formulation Taking Eddy Currents in Lamination Stacks Into Account,” *IEEE Trans. Magn.*, vol. 39, no. 3, pp. 1424-1427, May 2003.
- [4] L. Krähenbühl et al., “Homogenization of Lamination Stacks in Linear Magnetodynamics,” *IEEE Trans. Magn.*, vol. 40, no. 2, pp. 912 - 915 Mar. 2004.
- [5] H. De Gersem, S. Vanaverbeke, and G. Samaey, “Three-Dimensional-Two-Dimensional Coupled Model for Eddy Currents in Laminated Iron Cores,” *IEEE. Trans. Magn.*, vol. 48, no. 2, pp.815 – 818, Feb. 2012.